



This is Fig. 3*a* from Ryu & Jones (1995, ApJ, 442, 228), showing the solution of the MHD shock tube problem with the left state $(\rho, v_1, v_2, v_3, B_2, B_3, p_1) = [0.1, 50, 0, 0, -1/(4\pi)^{1/2}, -2/(4\pi)^{1/2}, 0.4]$ and the right state $[0.1, 0, 0, 0, 1/(4\pi)^{1/2}, 2/(4\pi)^{1/2}, 0.2]$ with $B_1 = 0$ and $\gamma = 5/3$ at time t = 0.01. At t = 0, the discontinuity is at $x_1 = 0.5$. Plots show from left to right: (1) magneto-acoustical shock, (2) tangential discontinuity (at $x_1 \sim 0.75$), and (3) magneto-acoustical shock. (When $B_1 = 0$, as in this case, neither slow nor Alfvén waves propagate, and the fast waves become magneto-acoustical waves.)

Open circles are the dzeus36 solution using 512 zones, CMoC, the total energy equation, and secondorder interpolation (no contact steepener). dzeus36 parameters controlling the time step and artificial viscosity are: courno=0.75, qcon=1.0, and qlin=0.2. Lines are the results from the non-linear Riemann solver described in Ryu & Jones.

There are no significant differences between the dzeus36 and dzeus35 solutions. The "glitch" in ρ , p_1 , and e_T at $x_1 \sim 0.75$ is numerical in origin, and appears in upwinded schemes such as TVD as well. It is a result of numerical and artificial viscous stresses at the original discontinuity, and can be nearly eliminated by setting the artificial viscous parameters to zero (at the expense of stability elsewhere).