

Huygens-Fresnel Principle

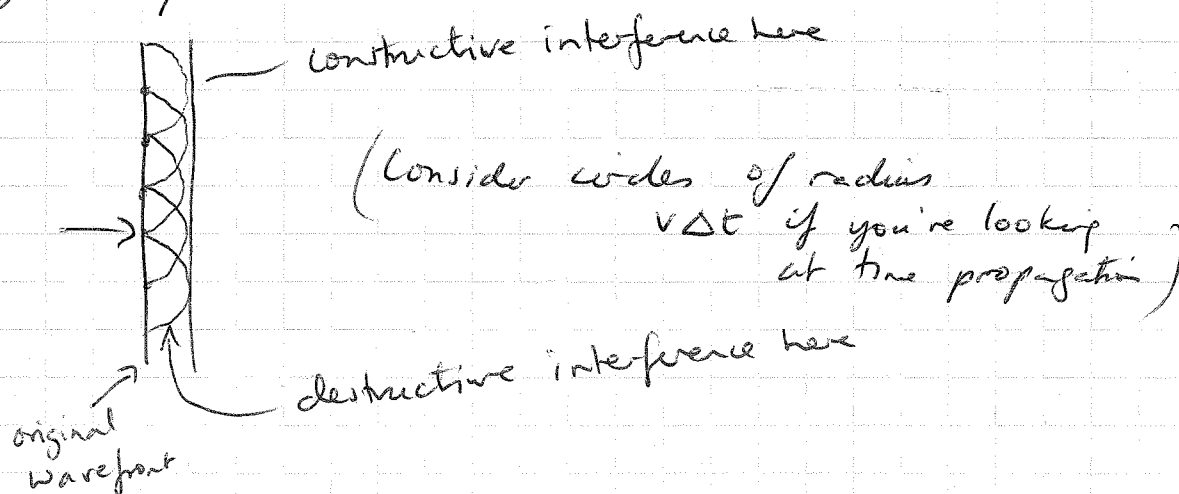
Q: How can we figure out how a wave will propagate through / around a material in 2D?

A: We need to look at how waves propagate - the Huygens-Fresnel Principle outlines a procedure that lets us do exactly this.

Key idea:

Each point in an advancing wave front is the source of new outgoing disturbance, with the resulting wave being a sum of all these outgoing disturbances.

i.e. for a plane wave:



Q: There is something slightly weird about this idea - can you see it?

A: The fact we only considered the forward moving wave - wouldn't there be one behind?

Actually we really need to consider the forward propagation only - after all a wave won't suddenly go backward without reflecting.

So we consider semi-circles in the direction of propagation.

This idea probably came from looking at water waves. If you put in a small gap in a sheet. Then the outgoing disturbance is a semi-circle regardless of the waveform hitting the gap. See Fig 8.6. "qualitative"

While this approach might feel slightly "artificial" it does give us a method of determining what happens when waves meet a barrier or obstacle.

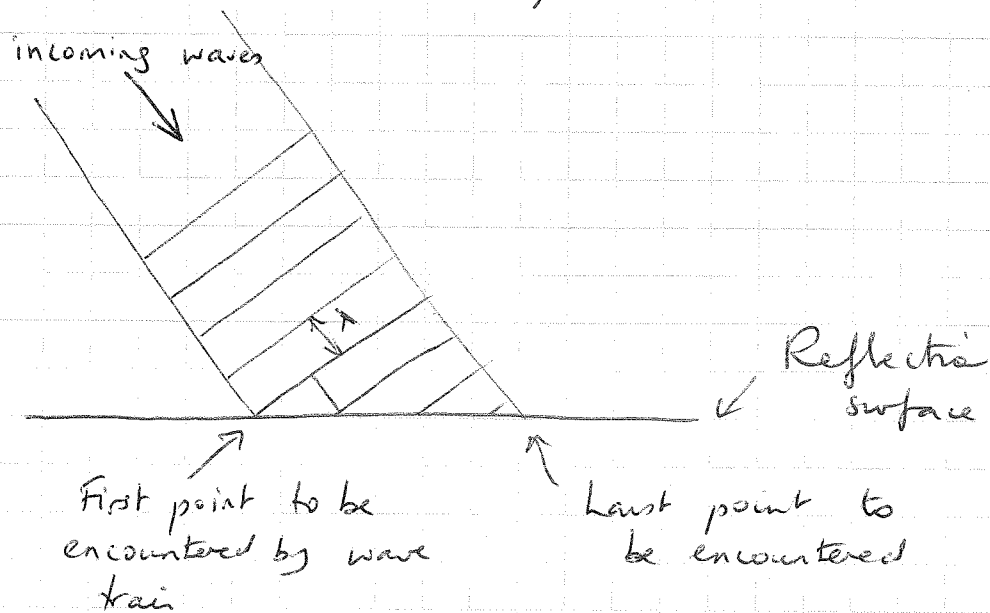
Reflection & Refraction of Plane Waves

We saw what happens in 1-d when a wave hits a discontinuity - reflection & partial transmission

Q: What can happen in 2-d that's different?

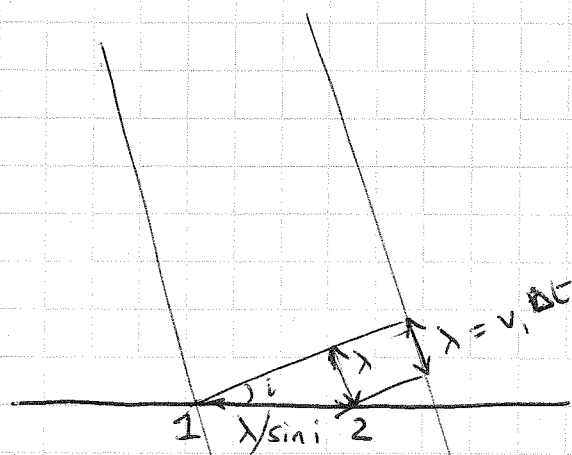
A: Change in the direction of propagation!

To analyze we'll just consider a train of incoming waves & we'll mark the peaks.



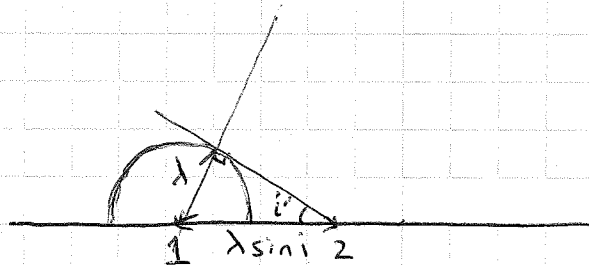
$$\text{Separation of wavefronts} = \lambda = v \Delta t$$

So now consider two (in phase) wavefronts & how far their disturbances will have propagated.



Now consider the wavefronts propagating from 1 & 2.

The wavefront 2 hasn't moved anywhere since it is the last of the two points. The wavefront from 1 will have moved $v_1 \Delta t = \lambda$



If we draw the tangent to the wavefront that will mark the front of the outgoing wave.

But the shape of this triangle is basically the same as what we had before

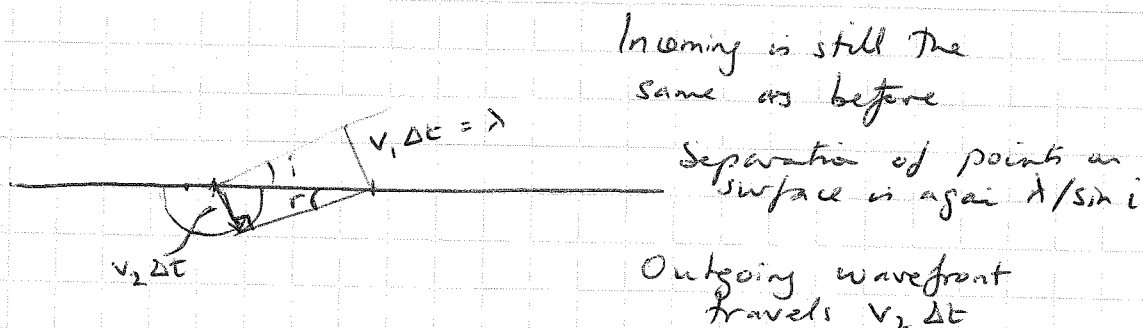
- opposite side length = λ
- hypotenuses " " = $\lambda / \sin i$

∴ $i = i'$ and the incoming angle of incidence matches the outgoing.

Exactly what we'd expect!

What about if we propagate into a different medium?

In this case the situation considers a forward moving wave rather than the reflected one. This medium will also have a different speed of propagation, say v_2



Incident is still the same as before

Separation of points on surface is again $\lambda/\sin i$

Outgoing wavefront travels $v_2 \Delta t$

How can we calculate angle r ?

It's again constructed as the tangent to the outgoing wavefront of radius $v_2 \Delta t$

The hypotenuse is again $\lambda/\sin i$ but using $\lambda = v_1 \Delta t$ we must have

$$\sin r = \frac{\text{opp}}{\text{hyp}} = \frac{v_2 \Delta t}{\lambda/\sin i} = \frac{v_2 \Delta t}{v_1 \Delta t/\sin i}$$

$$\Rightarrow \frac{\sin r}{\sin i} = \frac{v_2}{v_1}$$

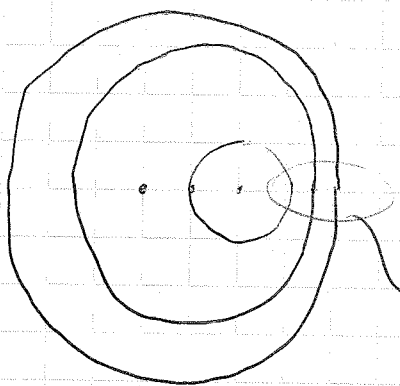
Note: evaluating the amplitudes is considerably harder. For transverse waves it even depends upon the state of polarization!

Doppler Effect & Related Phenomena

22.5

This is the situation where the source of the effect is moving through the medium, at say some speed u .

Case 1 $u < \text{sound speed}$



Consider wavefronts exactly one wavelength apart, i.e. time τ between them where $\lambda_0 = v\tau$

(τ is of course $\frac{1}{f}$)

BUT source moves $u\tau$ in the time between wavefronts.

wavefronts are compressed together

So separation between fronts is:

$$v\tau - u\tau = \lambda_{\text{new}}$$

$$\Rightarrow (v - u)\tau = \lambda_{\text{new}} \quad \& \text{ for situation at rest } \lambda_0 = v\tau$$

$$\Rightarrow \frac{(v - u)\tau}{v\tau} = \frac{\lambda_{\text{new}}}{\lambda_0}$$

$$\Rightarrow \lambda_{\text{new}} = \left(1 - \frac{u}{v}\right) \lambda_0 \quad \left(\text{Usually people just write } \lambda \text{ rather than } \lambda_{\text{new}}\right)$$

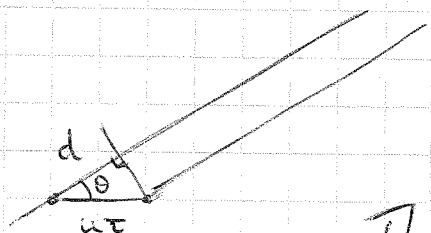
You can do the same analysis on the other side where the wavelength will be stretched out

$$\Rightarrow (v + u)\tau = \lambda_{\text{new}}$$

$$\Rightarrow \lambda_{\text{new}} = \left(1 + \frac{u}{v}\right) \lambda_0$$

So if you're behind the source you hear a longer wavelength, and ahead you hear a shorter wavelength.

What happens at an angle?



At rest we again have
 $\lambda_0 = v\tau$

The wave fronts propagate in circular directions as before

But net distance between wavefronts is
 $v\tau - d$
 $= v\tau - u\tau \cos \theta$

Quickly check this makes sense:

if straight ahead we have $v\tau - u\tau$ ✓
 if directly behind " " $v\tau + u\tau$ ✓

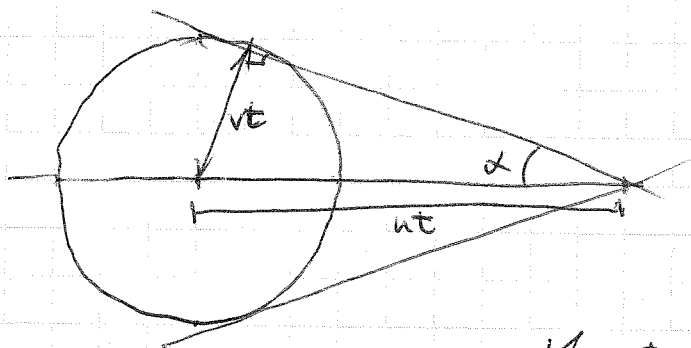
Thus $\lambda(\theta) = (v - u \cos \theta)\tau$

$$\Rightarrow \frac{\lambda(\theta)}{\lambda_0} = \frac{(v - u \cos \theta)\tau}{v\tau} = \left(1 - \frac{u}{v} \cos \theta\right)$$

$$\therefore \lambda(\theta) = \left(1 - \frac{u}{v} \cos \theta\right) \lambda_0$$

Situation 2: when $u > v$

Now the system moves faster than the sound speed.



$$\sin \alpha = \frac{v\tau}{u\tau} = \frac{v}{u}$$

Note $\frac{u}{v}$ = "Mach number"

α = "mach angle"

If these were sound waves this would be the "sonic boom" cone.