

# APPROXIMATIONS AT LARGE DEPTH,

$$\underline{T_v \gg 1:}$$

- LOWER BOUNDARY CONDITION

## DIFFUSION (ROSSELAND) APPROXIMATION:

AT SOME DEPTH,  $T_v$ :

$$S_v(T_v) \text{ \& \ } \left. \frac{d^n S_v(t_v)}{dt_v^n} \right|_{t_v=T_v} \text{ KNOWN, ALL } n$$

THEN, AT ANY DEPTH,  $t_v$ , TAYLOR

SERIES FOR  $S_v(t_v)$ :

$$S_v(t_v) = S(T_v) + \underbrace{(t_v - T_v)}_{\substack{\sim \\ \text{E-B}}} \left. \frac{dS_v(t_v)}{dt_v} \right|_{t_v=T_v} + \frac{(t_v - T_v)^2}{2} \left. \frac{d^2 S_v(t_v)}{dt_v^2} \right|_{t_v=T_v} + \dots$$

*2<sup>nd</sup> EDD.*

FORMAL Soln FOR  $I_\nu^+(\tau_\nu, \mu)$  &  $I_\nu^-(\tau_\nu, \mu)$ :

Eg.  $I_\nu^+(\tau_\nu, \mu) = \frac{1}{\mu} \int_{\tau_\nu}^{\infty} S_\nu(t_\nu) e^{-(t_\nu - \tau_\nu)/\mu} dt_\nu$

IN LIMIT OF  $\tau_\nu \rightarrow \infty$ , FOR  $-1 \leq \mu \leq 1$ :

$$I_\nu(\tau_\nu, \mu) \approx \underbrace{S_\nu(\tau_\nu)}_{2^{\text{nd}} \text{ EDD.}} + \overset{E-B}{\mu} \left. \frac{dS_\nu(t_\nu)}{dt_\nu} \right|_{\tau_\nu} + \mu^2 \left. \frac{d^2 S(t_\nu)}{dt_\nu^2} \right|_{\tau_\nu} + \dots$$

- RUTTEN, p. 87-88

- FOR  $I_\nu^+$  &  $I_\nu^-$

## ANGLE-MOMENTS:

$$J_v(\tau_v) \approx S_v(\tau_v) + \frac{1}{3} \frac{d^2 S_v(t_v)}{dt_v^2} \Big|_{\tau_v} + \dots$$

- EVEN TERMS

$$F_v(\tau_v) \approx \frac{4\pi}{3} \frac{dS_v(t_v)}{dt_v} \Big|_{\tau_v} + \frac{4\pi}{5} \frac{d^3 S_v(t_v)}{dt_v^3} \Big|_{\tau_v} + \dots$$

- ODD TERMS

$$K_v(\tau_v) \approx \frac{1}{3} S_v(\tau_v) + \frac{1}{5} \frac{d^2 S_v(t_v)}{dt_v^2} \Big|_{\tau_v} + \dots$$

- EVEN TERMS

AT  $\tau_v \rightarrow \infty$ ,  $S_v(\tau_v) \approx \underline{\mathcal{O}(1)}$ :

$$S_v(t_v) \approx S_v(\tau_v) + (t_v - \tau_v) \left. \frac{dS_v(t_v)}{dt_v} \right|_{\tau_v}$$

$\therefore \underline{\mathcal{O}(1)}$ :

$$I_v(\tau_v, \mu) \approx S_v(\tau_v) + \mu \left. \frac{dS_v(t_v)}{dt_v} \right|_{\tau_v}$$

$$J_v(\tau_v) \approx S_v(\tau_v) \text{ — (1)}$$

$$\tilde{F}_v(\tau_v) \approx \frac{4\pi}{3} \left. \frac{dS_v(t_v)}{dt_v} \right|_{\tau_v}$$

$$K_v(\tau_v) \approx \frac{1}{3} S_v(\tau_v) \text{ — (2)}$$

Eqs. (1) & (2) :

$$K_{\nu}(T_{\nu}) = \frac{1}{3} J_{\nu}(T_{\nu})$$

- 1<sup>ST</sup> EDDINGTON APPROX.

$$\therefore \frac{c}{4\pi} P_{\nu}(T_{\nu}) = \frac{1}{3} \frac{4\pi}{c} \mu_{\nu}(T_{\nu})$$

$$\therefore P_{\nu}(T_{\nu}) = \frac{1}{3} \mu_{\nu}(T_{\nu})$$

- SENSE

LTE:  $S_\nu(\tau_\nu) = B_\nu(T_{\text{KIN}}(\tau_\nu))$  :

$$I_\nu(\tau_\nu, \mu) \approx \underline{B}_\nu(\tau_\nu) + \mu \frac{dB_\nu(\tau_\nu)}{d\tau_\nu} \Big|_{\tau_\nu}$$

$$J_\nu(\tau_\nu) \approx \underline{B}_\nu(\tau_\nu) \quad - \text{SAME AS T.E.}$$

$$\underline{F}_\nu(\tau_\nu) \approx \frac{4\pi}{3} \frac{dB_\nu(\tau_\nu)}{d\tau_\nu}$$

RECALL:

IN 1D MODEL, RADIAL  $\tau_\nu$ -SCALE:

$$d\tau_\nu(\underline{z}) = -\alpha_\nu(\underline{z}) \underline{dz}$$

$$\underline{F}_\nu(\underline{z}) \approx \uparrow \frac{4\pi}{3} \frac{1}{\alpha_\nu(\underline{z})} \frac{dB_\nu(\underline{z})}{d\underline{z}}$$

- MONOCHROMATIC

## BOLOMETRIC FLUX, $F$ :

$$F(z) \approx -\frac{4\pi}{3} \int_{\nu=0}^{\infty} \frac{1}{\alpha_{\nu}(z)} \frac{dB_{\nu}(T(z))}{dz} d\nu$$

$$= -\frac{4\pi}{3} \left\langle \frac{1}{\alpha_{\nu}(z)} \right\rangle \frac{dT(z)}{dz} \int_{\nu=0}^{\infty} \frac{dB_{\nu}(T)}{dT} d\nu$$

## STEFAN-BOLTZMANN VARIATION:

$$\int_0^{\infty} \frac{B_{\nu}(T)}{dT} d\nu = \frac{4\sigma}{\pi} T^3$$

$$\therefore F(z) = -\frac{16\sigma}{\pi} \left\langle \frac{1}{\alpha_{\nu}(z)} \right\rangle T_{KIN}(z) \frac{dT_{KIN}(z)}{dz}$$

$$= -\beta(z) \frac{dT_{KIN}(z)}{dz}$$

- DIFFUSION Eq.

∴ DIFFUSION APPROX. ( $T_v \rightarrow \infty$ ):

$$\frac{dT_{KIN}(z)}{dz} = - \frac{F(z)}{\beta(z)}$$

- DIFFUSION CONSTANT:

$$\beta(z) = \frac{16\sigma}{\pi} \left\langle \frac{1}{\alpha_v(z)} \right\rangle T_{KIN}(z)$$

-  $\left\langle \frac{1}{\alpha_v} \right\rangle = V$ -AVERAGED MEAN  
EXTINCTION COEFFICIENT  
( $cm^{-1}$ )



## ROSSELAND MEAN EXTINCTION

COEFFICIENT,  $\alpha_R$  ( $\text{cm}^{-1}$ ):

$$\frac{1}{\alpha_R(z)} = \frac{\int_0^{\infty} \frac{1}{\alpha_\nu(z)} \frac{dB_\nu(z)}{dT} d\nu}{\int_0^{\infty} \frac{dB_\nu(z)}{dT} d\nu}$$

- HARMONIC MEAN

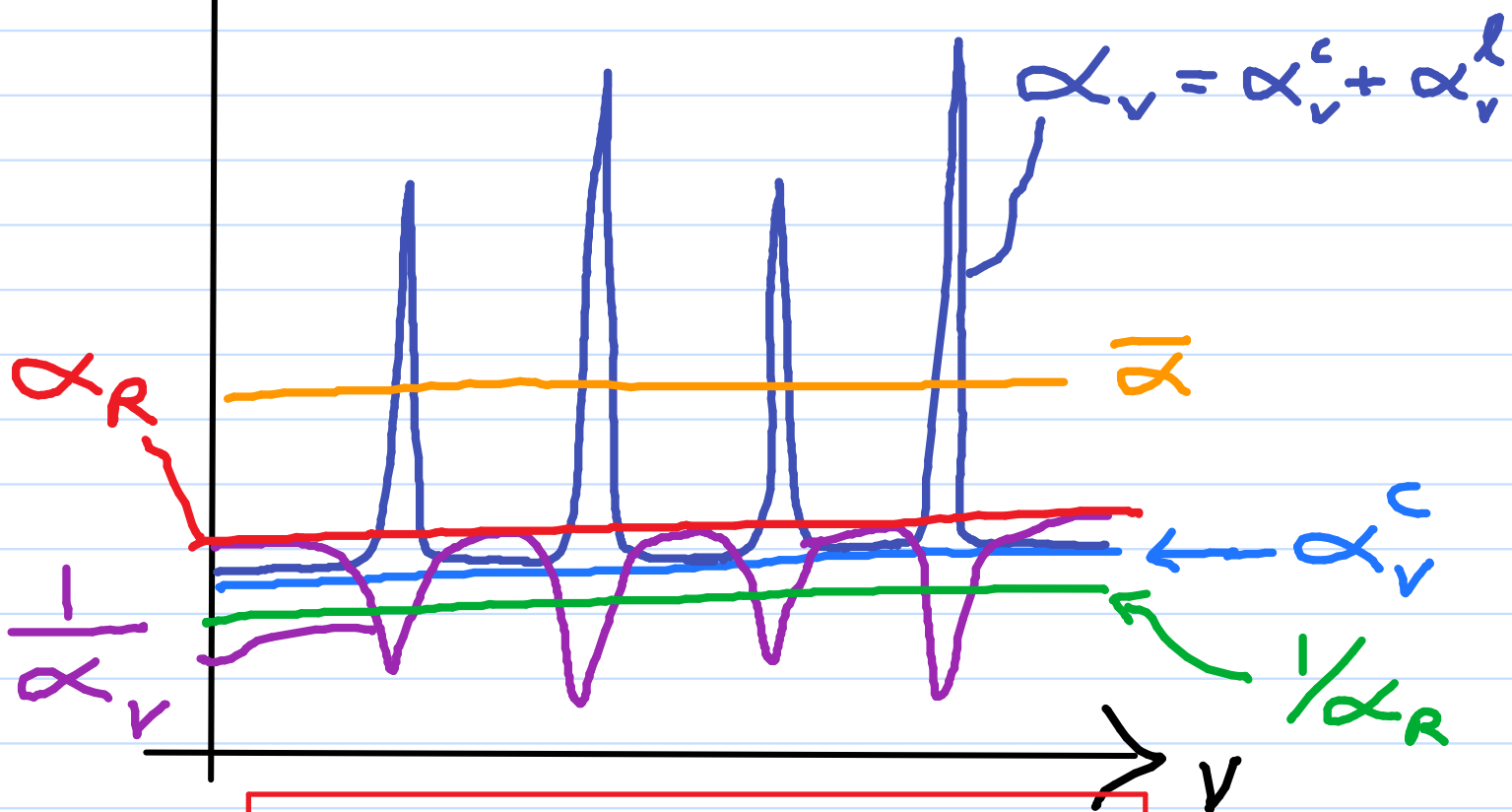
CONTRAST:

$$\text{STRAIGHT MEAN: } \bar{\alpha} = \frac{\int_0^{\infty} w_\nu \alpha_\nu d\nu}{\int_0^{\infty} w_\nu d\nu}$$

-  $w_\nu$  = WEIGHT fn

$\Delta\nu \approx \text{FEW} \times 10^{15} \text{ Hz}$

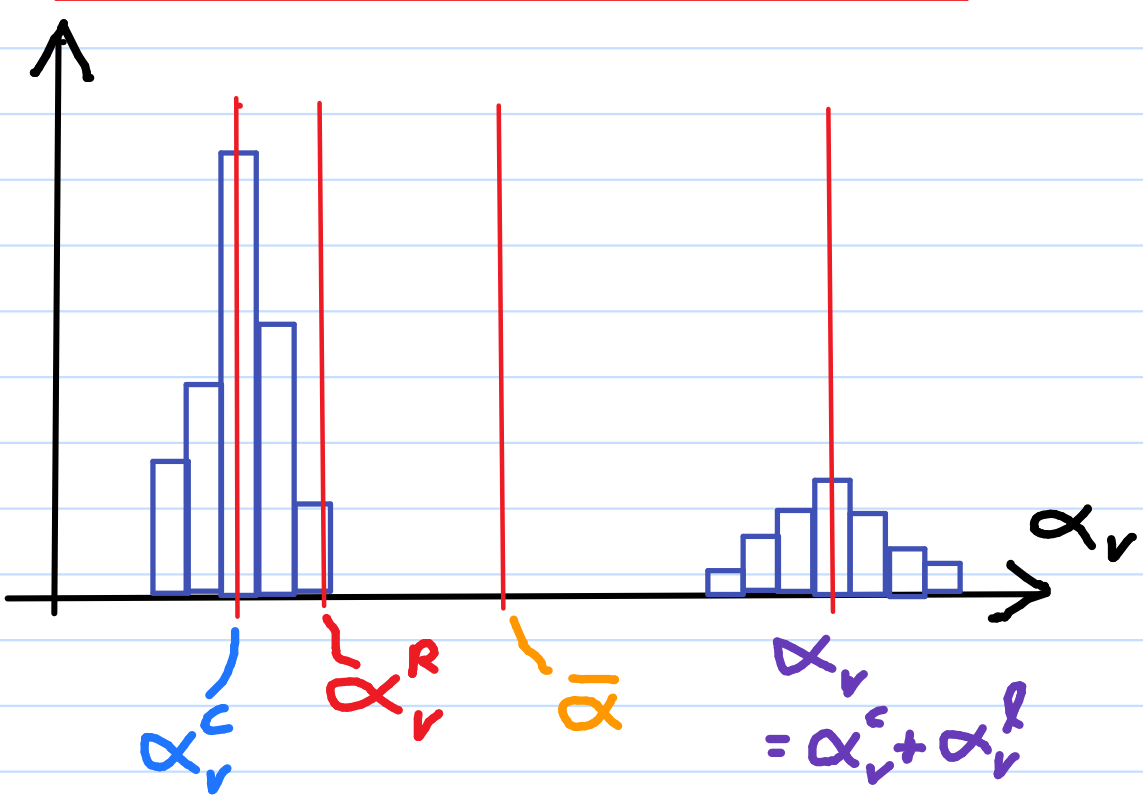
$\alpha_{\nu}$



$\alpha_{\nu} = \alpha_{\nu}^c + \alpha_{\nu}^l$

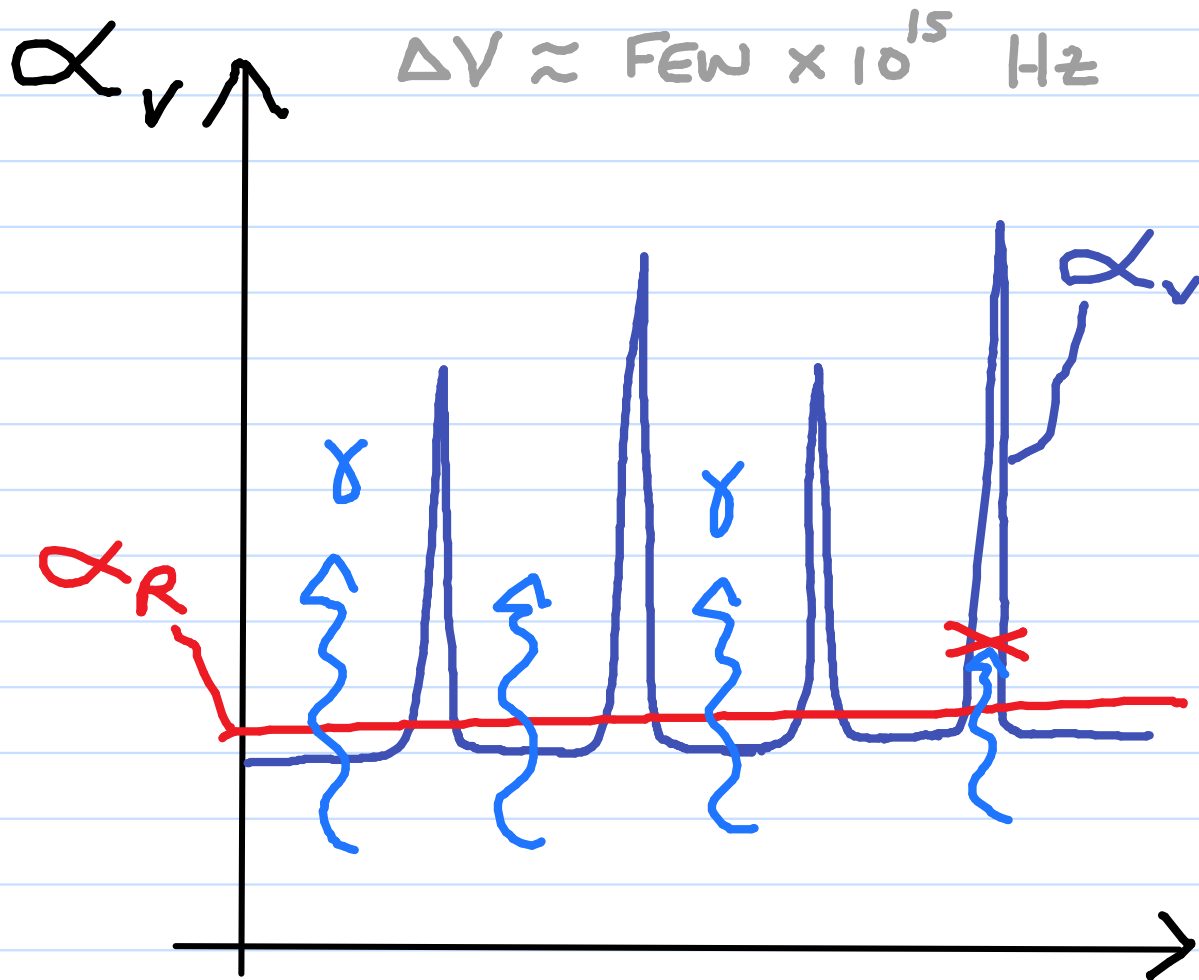
$\alpha_R \approx \alpha_{\nu}^c < \alpha$

$N$



$\alpha_{\nu}^c$     $\alpha_{\nu}^R$     $\alpha$     $\alpha_{\nu}^c + \alpha_{\nu}^l$

$\alpha_R$  IS A FLUX-WEIGHTED MEAN EXTINCTION



$$K_R(z) = \frac{\alpha_R(z)}{\rho} \quad (\text{cm}^2/\text{g})$$

$$dT_R(z) = -\alpha_R(z) dz$$

- ROSSELAND T-SCALE

THEN:

$$\kappa(z) \approx \frac{-16\sigma}{3} \frac{1}{\alpha_R(z)} T_{\text{KIN}}^3(z) \frac{dT_{\text{KIN}}(z)}{dz}$$

- IN LTE (ASSUMING  $T_{\text{RAD}}(z) = T_{\text{KIN}}(z)$ )

